Entanglement Assisted Covert Communication via Qubit Depolarizing Channels

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- Privacy and secrecy are critical in communication
- **Covert Communication**: Not only the transmitted information kept secret, but also the transmission itself.
- Information Theory: How many bits of information can be sent for *n* channel uses?
- Pre-shared entanglement resources can increase performance and throughput

- No-entanglement, $O(\sqrt{n})$ (SRL-square root law):
 - Classical communication [Bash et al. 2013, Bloch et al. 2016]
 - Discrete variable (classical-quatum) [Sheikholeslami et al. 2016]
 - Continuous variable (Gaussian bosonic) [Bash et al. 2015]
- Entanglement, $O(\sqrt{n}\log(n))$:
 - Continuous variable (Gaussian bosonic) [Gagatsos et al. 2020]
 - Discrete variable?



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- Entanglement, $O(\sqrt{n}\log(n))$:
 - Continuous variable (Gaussian bosonic) [Gagatsos et al. 2020]
 - Discrete variable? Yes!



- We consider covert communication over qubit depolarizing channels.
- Three scenarios:
 - Willie (adversary) has all environment
 - "half" the environment
 - the other "half"
- Logarithmic factor is not reserved for continuous variable
- Interpretation: energy-constrained capacities, decoding performance



Quantum Information: Pure States

A pure quantum state $|\psi\rangle$ is a normalized vector in the Hilbert space \mathcal{H}_A .

Qubit

For a quantum bit (qubit),

$$ert 0
angle = egin{pmatrix} 1 \ 0 \end{pmatrix} ert$$
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For a quantum bit (qubit),

$$\begin{split} |0\rangle &= \begin{pmatrix} 1\\ 0 \end{pmatrix} \\ |1\rangle &= \begin{pmatrix} \alpha\\ \beta \end{pmatrix} = \alpha |0\rangle + \beta |1\rangle \\ |1\rangle &= \begin{pmatrix} 0\\ 1 \end{pmatrix} \end{split}$$

 α and β can be complex numbers, and $|\alpha|^2+|\beta|^2=1$



A pure bi-partite state $|\psi_{AB}\rangle$ is a normalized vector in the product Hilbert space $\mathcal{H}_A \otimes \mathcal{H}_B$.

Two qubits

For two qubits, $|\psi_{AB}
angle = |i
angle \otimes |j
angle$, or

$$|\psi_{AB}
angle = \sum_{i,j=0,1} lpha_{ij} |i
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angle \ , \ ext{with} \ \sum |lpha_{ij}|^2 = 1$$



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Entanglement

Systems A and B are entangled if $|\psi_{AB}\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$

For example,
$$|\Phi_{AB}\rangle = \frac{1}{\sqrt{2}}(|0\rangle_A \otimes |0\rangle_B + |1\rangle_A \otimes |1\rangle_B).$$

The (mixed) state ρ_A of a quantum system A is an Hermitian, positive semidefinite, unit-trace density matrix over \mathcal{H}_A .

- Mixed states reflecting uncertainty about a quantum system.
- For a joint state ρ_{AB} , the density matrix of the system A is the reduced matrix $\rho_A = \text{Tr}_B(\rho_{AB})$.



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Entropy

Given ρ_A , define

$$H(A)_{\rho} \equiv -\mathrm{Tr}(\rho_A \log \rho_A)$$

For a **Pure** system *A*:

$$H(A)_{\rho} = 0$$

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Useful definitions:

• Divergence (Relative entropy):

 $D(
ho||\sigma) = \operatorname{Tr}[
ho \log(
ho) -
ho \log(\sigma)]$

(if $\operatorname{supp}(\rho) \subseteq \operatorname{supp}(\sigma)$; and $D(\rho||\sigma) = +\infty$, otherwise.)



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• Second moment:

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• η -divergence: For a spectral decomposition $\sigma = \sum_i \lambda_i P_i$, let [Tahmasbi et al. 2021]:

• Pauli Matrices:

$$X = \begin{pmatrix} 0 & 1 \\ 1 & 0 \end{pmatrix}, \quad Y = \begin{pmatrix} 0 & -i \\ i & 0 \end{pmatrix}, \quad Z = \begin{pmatrix} 1 & 0 \\ 0 & -1 \end{pmatrix}$$



Quantum Channel

Unitary Evolution - Pure State

• Evolution of a pure state is given by a unitary operator in Hilbert space \mathcal{H}_A :

$$|\psi
angle \stackrel{U}{\longrightarrow} U|\psi
angle \qquad U^{\dagger}U = UU^{\dagger} = \mathbb{1}$$



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Noisy Evolution - Density Matrix

- Evolution of a density matrix is given by a noisy channel $\mathcal{N}_{A \rightarrow B}$.
- A quantum noisy channel is defined as a completely-positive trace-preserving (CPTP) linear map.

$$\rho_A \xrightarrow{\mathcal{N}} \rho_B \equiv \operatorname{Tr}_E(V \rho_A V^{\dagger}) \qquad \qquad V \equiv V_{A \to BE}^{\mathcal{N}}$$
$$V^{\dagger} V = \mathbb{1}_A$$

• $V_{A \to BE}^{\mathcal{N}}$ is a linear opertor from \mathcal{H}_A (Alice) to $\mathcal{H}_B \otimes \mathcal{H}_E$ (Bob and Enviroment).

Example: "qubit flip" channel.

 \bullet Flips $|0\rangle$ into $|1\rangle$ and $|1\rangle$ into $|0\rangle$ with probability of q:

$$\mathcal{N}_{A
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ho) = (1-q)
ho + qX
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• The "qubit flip" channel can be given by:

$$\mathcal{N}_{A \to B}(\rho) = \operatorname{Tr}_{E}(V \rho V^{\dagger})$$

where $V : \mathcal{H}_A \to \mathcal{H}_B \otimes \mathcal{H}_E$ is an isometry defined by

$$V_{A
ightarrow BE} \equiv \sqrt{1-q} \mathbb{1} \otimes \ket{1} + \sqrt{q} X \otimes \ket{2} \, .$$



Quantum Channel: Example 2 (Study Case)

Example: qubit depolarizing channel.

 Bob receives Alice's qubit with probability 1 − q, and a completly mixed state with probability q:

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Without Covertness



- $\log(M)$ #information bits, over *n* channel uses.
- Rate: $R = \frac{\log(M)}{n}$
- Example: 3-repetition code. 0 ightarrow |000angle, 1 ightarrow |111angle
 - $n = 3 \cdot \log(M)$, hence: $R = \frac{1}{3}$
 - #information bits: $\log(M) = \frac{1}{3} \cdot n = O(n)$
- In covert communication, log(M) is sub-linear



Covert Communication



• Reliability: Bob's probability of error tends to zero

 $\lim_{n\to\infty} \Pr(\text{error}) = 0$

• Covertness: Willie has a bad detection performance

$$\lim_{n\to\infty} D(\bar{\rho}_{W^n}||\omega_0^{\otimes n}) = 0$$

• Covert "rate":

$$L = \frac{\log(M)}{\log(n)\sqrt{nD(\overline{\rho}_{W^n}||\omega_0^{\otimes n})}}$$

.

• Covert capacity: Supremum achievable rate as $n \to \infty$

- The scale of O(√n log(n)) has already been shown in a continuous-variable model, i.e., the bosonic Gaussian channel [Gagatsos et al. 2020].
- Until now, it has remained unclear whether this performance boost can also be achieved in finite dimensions.
- There are communication settings, which the coding scale is larger for continuous-variable channels.
- For example, in deterministic identification, the code size is super-exponential for Gaussian channels but limited to an exponential scale for finite-dimensional channels [Salariseddigh et al. 2021].



Depolarizing Channel

- We study communication over the depolarizing channel.
- The qubit depolarizing channel can be given by:

$$\mathcal{N}_{A\to B}(\rho) = \mathrm{Tr}_{E_1E_2}(V\rho V^{\dagger})$$

where $V:\mathcal{H}_A\to\mathcal{H}_B\otimes\mathcal{H}_{E_1}\otimes\mathcal{H}_{E_2}$ is an isometry defined by

$$V\equiv \sqrt{1-rac{3q}{4}}\mathbbm{1}\otimes \ket{00} + \sqrt{rac{q}{4}}X\otimes \ket{01} + \sqrt{rac{q}{4}}Y\otimes \ket{11} + \sqrt{rac{q}{4}}Z\otimes \ket{10} \,.$$

- Three qubits at the output of the channel
- 1st qubit belongs to Bob. 2nd and 3rd leak to the environment.
- Intuitively, (E_1, E_2) store a "flag" that indicates which Pauli error occurred.
- Willie's access is limited to the environment.
- No-Cloning Theorem: Willy's channel cannot be the same as Bob's

Willie's Channel

Willie has an access to (part of) the environment. We consider three cases:

- Scenario 1: Willie receives both qubits, E_1 and E_2 .
- Scenario 2: Willie receives last qubit, E_2 .

• Scenario 3: Willie receives the qubit E_1 .



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Willie's Channel: Scenario 1

Theorem

Covert communication is impossible in Scenario 1. Hence, if $W = (E_1, E_2)$, then $C_{cov-EA}(\mathcal{N}) = 0$.

- Willie receives the entire environment
- This is strong enough for him to detect any encoding operation.
- $\operatorname{supp}(\omega_1) \not\subseteq \operatorname{supp}(\omega_0)$, where $\omega_0 \equiv \widehat{\mathcal{N}}_{A \to W}(|0\rangle\langle 0|)$ and $\omega_1 \equiv \widehat{\mathcal{N}}_{A \to W}(|1\rangle\langle 1|)$
- Note: ω_1 and ω_0 depend on the channel parameter q.



Willie's Channel: Scenario 2

Theorem

Covert communication is trivial in Scenario 2. That is, Alice can communicate information as without the covertness requirement, and send O(n) bits.

- Willie receives the second qubit.
- \bullet Willie cannot discern between the $|0\rangle$ and $|1\rangle$ inputs.
- $\omega_0 = \omega_1 = \left(1 \frac{q}{2}\right) \left|0\right\rangle \left(0\right| + \frac{q}{2} \left|1\right\rangle \left(1\right|$



Willie's Channel: Scenario 3

- Willie receives the first qubit.
- Covert communication is possible, and not trivial
- $\operatorname{supp}(\omega_1) \subseteq \operatorname{supp}(\omega_0)$ and $\omega_0 \neq \omega_1$



Theorem

Consider a qubit depolarizing channel as in scenario 3. The entanglement-assisted covert capacity is bounded as

$$\mathcal{C}_{cov\text{-}\mathcal{E}\mathcal{A}}(\mathcal{N}) \geq rac{4\sqrt{2}}{3}rac{(1-q)^2}{(2-q)\sqrt{\eta(\omega_1||\omega_0)}}$$

where $\omega_0 \equiv \mathcal{N}_{A \to W}(|0\rangle\!\langle 0|)$ and $\omega_1 \equiv \mathcal{N}_{A \to W}(|1\rangle\!\langle 1|)$.

- Reminder covert capacity is the supremum of $\frac{\log(M)}{\log(n)\sqrt{nD(\overline{\rho}_{W^n}||\omega_n^{\otimes n})}}$
- Without entanglement, #information bits follows SRL, and here, the rate is defined according to the $\sqrt{n}\log(n)$ scale.
- Covert transmission of $O(\sqrt{n} \log n)$ information bits is achievable.
- Entanglement leads to a logarithmic performance boost.

Main Results: Lower Bound

Lower bound of the covert rate $C_{\text{cov-EA}}$ as function of the noise parameter q:



- $q \rightarrow 0$: No noise, covert communication is trivial.
- $q \rightarrow 1$: Completely noise, communication is impossible.

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Lemma (Wilde 2017, Gagatsos et al. 2020)

For an input state ψ_{A_1A} , and sufficiently large n, there exists a coding scheme that employs pre-shared entanglement resources to transmit $\log(M)$ bits over n uses of $\mathcal{N}_{A \rightarrow B}$ such that:

$$\log(M) \geq nD(\psi_{A_{1}B}||\psi_{A_{1}} \otimes \psi_{B}) + \sqrt{nV(\psi_{A_{1}B}||\psi_{A_{1}} \otimes \psi_{B})}\Phi^{-1}(\varepsilon) - C_{n}$$

with

$$\psi_{A_{\mathbf{1}}B} = (\mathrm{id}_{A_{\mathbf{1}}} \otimes \mathcal{N}_{A \to B})(\psi_{A_{\mathbf{1}}A})$$

where,

$$C_n = \frac{\beta_{\mathsf{B}-\mathsf{E}}}{\sqrt{2\pi}} \frac{[Q(\psi_{\mathsf{A}_1\mathsf{B}}||\psi_{\mathsf{A}_1}\otimes\psi_{\mathsf{B}})]^{\frac{3}{4}}}{V(\psi_{\mathsf{A}_1\mathsf{B}}||\psi_{\mathsf{A}_1}\otimes\psi_{\mathsf{B}})} + \frac{V(\psi_{\mathsf{A}_1\mathsf{B}}||\psi_{\mathsf{A}_1}\otimes\psi_{\mathsf{B}})}{\sqrt{2\pi}} + \log(4\varepsilon n),$$

and Φ^{-1} is the inverse-Gaussian distribution function.

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with

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- The derivation is based on a *position-based* coding scheme.
- Each message is associated with *n* entangled pairs
- Bob uses sequential decoding on the output and the entanglement resources for each message consecutively

- Unassisted communication: classical encoding [Sheikholeslami et al. 2016]
 - Alice selects binary sequences according to Bernoulli(α_n), where $\alpha_n \sim \frac{1}{\sqrt{n}}$.
 - The average input state is $\psi_A = (1 \alpha_n) |0\rangle\langle 0| + \alpha_n |1\rangle\langle 1|$.



- Unassisted communication: classical encoding [Sheikholeslami et al. 2016]
 - Alice selects binary sequences according to Bernoulli(α_n), where $\alpha_n \sim \frac{1}{\sqrt{n}}$.
 - The average input state is $\psi_A = (1 \alpha_n) |0\rangle \langle 0| + \alpha_n |1\rangle \langle 1|$.
- Entanglement-assisted communication: In our scheme,
 - Alice encodes with the superposition state:

$$|\psi_{A_1A}\rangle = \sqrt{1 - \alpha_n} |00\rangle + \sqrt{\alpha_n} |11\rangle ,$$

where $\alpha_n \sim \frac{1}{\sqrt{n}}$.

- A_1 is the pre-shared entanglement resource of Bob.
- $|\psi_{A_1A}\rangle$ can be considered as "very close" to the innocent state $|00\rangle$.
- The channel input A is the reduced state $\psi_A = (1 \alpha_n) |0\rangle\langle 0| + \alpha_n |1\rangle\langle 1|$.

Alice encodes with the state:

$$|\psi_{A_1A}\rangle = \sqrt{1 - \alpha_n} |00\rangle + \sqrt{\alpha_n} |11\rangle ,$$

using the mentioned lemma, and after some algebraic manipulation, we obtain

$$\log(M) \geq -2\frac{(1-q)^2}{2-q}\alpha_n\log(\alpha_n) + O(\alpha_n),$$

and finally

$$\mathcal{C}_{\mathsf{cov-EA}}(\mathcal{N}) \geq rac{4\sqrt{2}}{3} rac{(1-q)^2}{(2-q)\sqrt{\eta(\omega_1||\omega_0)}}$$



Main Result: Info. Bits Graph

Number of information bits for noise parameter $q = \frac{1}{2}$, and $D(\bar{\rho}_{W^n} || \omega_0^{\otimes n}) \leq 0.1$:





- The total energy of a state must not exceed a certain limit.
- A state ρ satisfies the energy constraint E, with the Hamiltonian \hat{H} , if $Tr(\hat{H}\rho) \leq E$.
- For $E \ll 1$,
 - $\bullet~$ Unassisted energy-constrained capacity: $\mathit{C}_0 \sim \mathit{E}$
 - $\bullet~$ Entanglement assisted energy-constrained capacity: $\mathit{C_{EA}}\sim -\mathit{E}\log\mathit{E}$
- The ratio between the assisted and unassisted scales as $-\log(E)$
- For $E_n \sim \frac{1}{\sqrt{n}}$, the ratio scales as $\log(n)$
- Effectively, the covertness requirement imposes an energy constraint. with Hamiltonian $\hat{H} = |1\rangle\langle 1|$ and the constraint $E_n \sim \frac{1}{\sqrt{n}}$



The "unfair channel setting": Bob can determine that some outputs are associated with a non-zero input, while Willie cannot. Hence, Bob has an unfair advantage over Willie.

- Examples: erasure channel, amplitude-damping channel.
- Even without assistance, # information bits scales as $\sqrt{n}\log(n)$ [Bloch et al. 2016, Sheikholeslami et al. 2016]

The depolarizing channel is fair in this sense, yet entanglement assistance has a similar effect as granting Bob the capability of identifying a non-zero transmission with certainty.



- We address covert communication over depolarizing channels
 - * main question: how can entanglement resources improve performance?
- We consider three scenarios: Willie has the entire environment, or, part of it.
- Our main contributions include:
 - $\ast\,$ Analysis of # information bits.
 - * Demonstrating that the logarithmic factor is not exclusive to continuous variable systems.
 - * Interpretation of covert communication rates as energy-constrained capacities for the qubit depolarizing channel.



- Entanglement assisted covert communication over a general channel
- Quantum covert communication
- Converse upper bound



Thank you



Appendix A - Divergence of two Bernoulli's



• Divergence between Bernoulli(x) and (0.5) $D(||0,5) = b_{1,2}(x, y) + (1, -x) b_{2,2}(1-x)$

•
$$D(x||0.5) = x \log(\frac{x}{0.5}) + (1-x) \log(\frac{1-x}{0.5})$$



Appendix B - explicit expression of ω_0 and ω_1 - scenario 1

$$\omega_{0}(q) = \begin{pmatrix} 1 - \frac{3q}{4} & 0 & 0 & \sqrt{\frac{q}{4}(1 - \frac{3q}{4})} \\ 0 & \frac{q}{4} & -i\frac{q}{4} & 0 \\ 0 & i\frac{q}{4} & \frac{q}{4} & 0 \\ \sqrt{\frac{q}{4}(1 - \frac{3q}{4})} & 0 & 0 & \frac{q}{4} \end{pmatrix}$$
(1)
$$\omega_{1}(q) = \begin{pmatrix} 1 - \frac{3q}{4} & 0 & 0 & -\sqrt{\frac{q}{4}(1 - \frac{3q}{4})} \\ 0 & \frac{q}{4} & i\frac{q}{4} & 0 \\ 0 & -i\frac{q}{4} & \frac{q}{4} & 0 \\ -\sqrt{\frac{q}{4}(1 - \frac{3q}{4})} & 0 & 0 & \frac{q}{4} \end{pmatrix}$$
(2)

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Appendix C - explicit expression of $\mathcal{N}_{A \to W}(\rho)$ - scenario 2

$$\mathcal{N}_{A \to W}(\rho) = \left(1 - \frac{q}{2}\right) |0\rangle\langle 0| + \frac{q}{2} |1\rangle\langle 1| + 2\operatorname{Re}\left\{b\right\} \left(\left(\sqrt{\left(1 - \frac{3q}{4}\right)\frac{q}{4}} + i\frac{q}{4}\right) |0\rangle\langle 1| + \left(\sqrt{\left(1 - \frac{3q}{4}\right)\frac{q}{4}} - i\frac{q}{4}\right) |1\rangle\langle 0|\right)$$
(3)



• Random codebook generation: Randomly and independently generate $2^{\log(M)}$ sequences (codewords) $x^n(m)$, $m \in [1 : M]$

$$x^{n}(m) \sim \prod_{i=1}^{n} p_{X}(x_{i})$$
(4)

- **Encoding:** To send message $m \in [1 : M]$, use $x^n(m)$.
- Decoding: Given a received sequence yⁿ, the decoder searches for a codeword xⁿ in the set of possible transmitted codewords such that (xⁿ, yⁿ) are jointly typical. (vary close to the expected probability given by P_{XY}).
- Why does it work? Law of large numbers.



Appendix E - Hypothesis testing relative entropy

• The hypothesis testing relative entropy is defined for $\varepsilon \in [0,1]$ as :

$$\mathcal{D}_{\mathcal{H}}^{\varepsilon}(\rho||\sigma) = -\log\inf_{\Lambda} \{ \operatorname{Tr}\{\Lambda\sigma\} : \operatorname{Tr}\{\Lambda\rho\} \ge 1 - \varepsilon \land 0 \le \Lambda \le I \} \,. \tag{5}$$

• The following expansion holds for a sufficiently large positive integer n:

$$D_{H}^{\varepsilon}(\rho^{\otimes}||\sigma^{\otimes}) = nD(\rho||\sigma) + \sqrt{nV(\rho||\sigma)}\Phi^{-1}(\varepsilon) + O(\log n)$$
(6)

where:

$$\Phi^{-1}(\varepsilon) = \sup\{\varepsilon \in \mathbb{R} | \Phi(\varepsilon) \le \varepsilon\} \qquad \Phi(\varepsilon) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\varepsilon} dx \exp\left(-x^2/2\right)$$
(7)

 Φ(ε) comes from Berry–Esseen theorem - a variation of the central limit theorem.