

Quantum Channel State Masking

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Quantum Communication and Information Theory

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 - Quantum key distribution for secure communication (307 km in optical fibers, 1200 km through space)
 - Computation power: Google’s supremacy experiment

State-dependent channels

- Channel state information (CSI)
 - classical applications: cognitive radio in wireless systems, memory storage, digital watermarking, etc.
- **State masking**: the state sequence represents information that should remain hidden from the receiver [Merhav and Shamai, 2007]

Classical results with channel state information (CSI) at the encoder:

- Causal CSI [Shannon 1958]
- Strictly-causal CSI [Csiszár and Körner 1981]
- Non-causal CSI [Gel'fand and Pinsker 1980]

Background: State-Dependent Channels

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Classical [state masking](#) [Merhav and Shamai, 2007]

- Broadcast channel [Koyluoglu et al. 2016] [Dikshitein et al. 2019]
- Source coding [Courtade 2012]
- Coordination [Le Treust and Bloch 2020]

Quantum channels with side information

- Without entanglement assistance: classical-quantum channels with causal or non-causal CSI [Boche, Cai, and Nötzel 2016]
- Entanglement assistance with non-causal CSI [Dupuis 2008]
- Entanglement assistance with causal CSI [P. 2020]
- Rate & State channel (parameter estimation) [P. 2021]

- Rate-limited entanglement assistance
 - Achievable rate-leakage region: tradeoff between communication, leakage, and entanglement resources.

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- Proof:
 - Achievability is based on the decoupling approach
 - Converse proof: classical arguments do not work

- Definitions
- Main Results
- Example
- Concluding Remarks

A pure quantum state $|\psi\rangle$ is a normalized vector in the Hilbert space \mathcal{H}_A .

Qubit

For a qubit, $|\psi\rangle = |0\rangle, |1\rangle$, or

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Quantum States

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Entanglement

Systems A and B are entangled if $|\psi_{AB}\rangle \neq |\psi_A\rangle \otimes |\psi_B\rangle$

For example, $|\Phi_{AB}\rangle = \frac{1}{\sqrt{2}}(|0\rangle \otimes |0\rangle + |1\rangle \otimes |1\rangle)$.

Entanglement can generate shared randomness, but it is a much more powerful resource.

Quantum States (Cont.)

The state ρ_A of a quantum system A is an Hermitian, positive semidefinite, unit-trace density matrix over \mathcal{H}_A .

Given ρ_{AB} , define

$$H(A)_\rho \equiv -\text{Tr}(\rho_A \log \rho_A)$$

$$H(A|B)_\rho \equiv H(AB)_\rho - H(B)_\rho$$

Quantum States (Cont.)

- Mutual information $I(A; B)_\rho = H(A)_\rho + H(B)_\rho - H(AB)_\rho$
- Coherent information $I(A \rangle B)_\rho = -H(A|B)_\rho$.

Quantum state-dependent channel

- A given CPTP linear map $\mathcal{N}_{EA \rightarrow B}$
- A pure state $|\phi_{EE_0C}\rangle^{\otimes n}$ (memoryless)
- Channel state information (CSI): Alice has E_0^n
- Entanglement resources: Alice and Bob share $\Psi_{G_A G_B}$

Quantum state-dependent channel

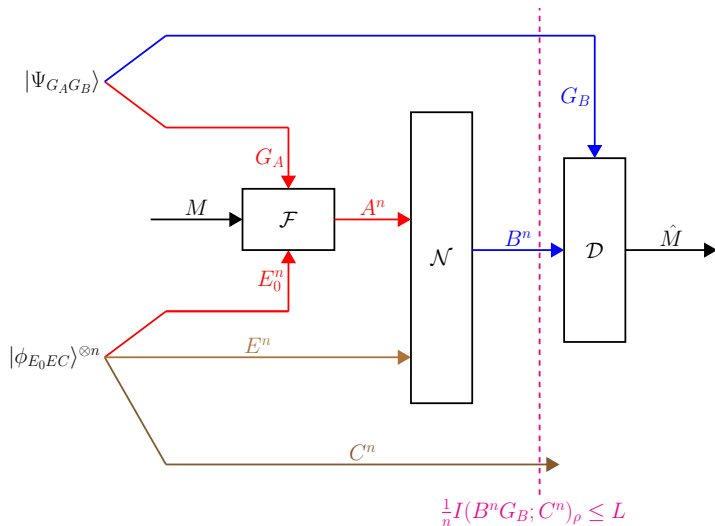
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Leakage

The state of C^n should be hidden from Bob

E.g., network information that should not be leaked to the end user.

Channel Model (Cont.)



Leakage

- In the classical case, the leakage requirement need not include shared randomness (cannot help the decoder).
- Our leakage requirement includes the entanglement resource system because Bob could use it to extract information on the channel state, using teleportation for instance.

Outline

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Main Results: Rate-Limited Assistance

Theorem

A quantum communication rate Q is achievable with leakage rate L and entanglement-assistance rate R_e if

$$Q + R_e \leq H(A|EC)_\rho$$

$$Q - R_e \leq -H(A|B)_\rho$$

$$L \geq I(C; AB)_\rho$$

for some $\rho_{AA'EC}$ with $\rho_{EC} = \phi_{EC}$, where $\rho_{ABC} = \mathcal{N}_{EA' \rightarrow B}(\rho_{AEA'C})$.

- demonstrates tradeoff between communication, leakage, and entanglement rates.
- Proof is based on the decoupling approach.

Proof

Main Results: Rate-Limited Assistance (Cont.)

- Masking is a “decoupling problem”: We wish to decouple C^n from B^n and G_B .
- In our extended decoupling approach, Bob's environment + C^n are decoupled from Alice's reference system.
- The leakage derivation follows naturally.

Main Results: Unassisted Region

Define

$$\underline{Q}(\mathcal{N}) \equiv \bigcup_{\rho_{EA'AC} : \rho_{EC} = \phi_{EC}} \left\{ (Q, L) : \begin{array}{l} 0 \leq Q \leq \min\{-H(A|B)_\rho, H(A|EC)_\rho\} \\ L \geq I(C; AB)_\rho \end{array} \right\}$$

and

$$\overline{Q}(\mathcal{U}^H) \equiv \bigcup_{\rho_{EA'AC} : \rho_{EC} = \phi_{EC}} \left\{ (Q, L) : \begin{array}{l} 0 \leq Q \leq H(A|CK)_\rho \\ L \geq I(C; AB)_\rho \end{array} \right\}$$

with $\rho_{ABC} = \mathcal{N}_{EA' \rightarrow B}(\rho_{AEA'C})$ and $\rho_{ABKC} = \mathcal{U}_{EA' \rightarrow BK}^H(\rho_{AEA'C})$.

Theorem

- 1) *the quantum masking region without assistance is given by*

$$\mathcal{R}_Q = \bigcup_{k=1}^{\infty} \frac{1}{k} \underline{Q}(\mathcal{N}^{\otimes k}).$$

- 2) *For a Hadamard channel,*

$$\underline{Q}(\mathcal{N}^H) \subseteq \mathcal{R}_Q \subseteq \overline{Q}(\mathcal{U}^H)$$

- arguments of Merhav and Shamai (2007) do **not** work in the quantum setting because $H(M|B^n C^n)_\rho < 0$.

Converse Proof (Unassisted)

The leakage is bounded by

$$\begin{aligned}n(L + \delta_n) &\geq I(C^n; B^n)_\rho \\ &= I(C^n; MB^n)_\rho - I(C^n; M|B^n)_\rho \\ &= I(C^n; MB^n)_\rho - H(M|B^n)_\rho + H(M|B^n C^n)_\rho\end{aligned}$$

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Since $H(M|B^n C^n)_\rho \geq -\log |\mathcal{H}_M| = -nQ$,

$$L + \delta_n + \varepsilon_n \geq \frac{1}{n} I(C^n; MB^n)_\rho$$

Main Results: Entanglement-Assisted Region

Theorem

Given entanglement assistance, the quantum capacity-leakage region is

$$\mathcal{R}_Q^{ea} = \bigcup_{\rho_{EA'AC} : \rho_{EC} = \varphi_{EC}} \left\{ (Q, L) : \begin{array}{l} 0 \leq Q \leq \frac{1}{2}[I(A; B)_\rho - I(A; EC)_\rho] \\ L \geq I(C; AB)_\rho \end{array} \right\}$$

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and the classical capacity-leakage region is

$$\mathcal{R}_{Cl}^{ea} = \bigcup_{\rho_{EA'AC} : \rho_{EC} = \varphi_{EC}} \left\{ (R, L) : \begin{array}{l} 0 \leq R \leq I(A; B)_\rho - I(A; EC)_\rho \\ L \geq I(C; AB)_\rho \end{array} \right\}$$

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assuming **maximally correlated** channel state systems:

$$\varphi_{EE_0C} = \sum_{s \in S} q(s) |s\rangle\langle s|_E \otimes |s\rangle\langle s|_{E_0} \otimes |s\rangle\langle s|_C$$

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Example: Dephasing Channel

State-dependent dephasing channel

Given a classical channel state $S \sim \text{Bernoulli}(q)$,

$$\mathcal{N}_{EA \rightarrow B}(\rho_{EA}) = (1 - q)\mathcal{P}_{A \rightarrow B}^{(0)}(\sigma_0) + q\mathcal{P}_{A \rightarrow B}^{(1)}(\sigma_1)$$

$$\mathcal{P}_{A \rightarrow B}^{(s)}(\sigma) = (1 - \varepsilon_s)\sigma + \varepsilon_s Z \sigma Z, \quad s = 0, 1,$$

for $\rho_{EA} = (1 - q)|0\rangle\langle 0|_E \otimes \sigma_0 + q|1\rangle\langle 1|_E \otimes \sigma_1$.

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If Alice applies Z gate controlled by $S \oplus Y$, $Y \sim \text{Bernoulli}(\lambda)$, this achieves

$$\mathcal{R}_{\text{Cl}}^{\text{ea}} \supseteq \bigcup_{0 \leq \lambda \leq \frac{1}{2}} \left\{ \begin{array}{l} (R, L) : 0 \leq R \leq 2 - h_2(\lambda * \hat{\varepsilon}) \\ L \geq h_2(\lambda * \hat{\varepsilon}) - (1 - q)h_2(\lambda * \varepsilon_0) - qh_2(\lambda * \varepsilon_1) \end{array} \right\}$$

where $a * b = (1 - a)b + a(1 - b)$ and $\hat{\varepsilon} = (1 - q)\varepsilon_0 + q(1 - \varepsilon_1)$. 

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Concluding Remarks

Our results demonstrate the following common phenomena in quantum information theory:

- Entanglement-assisted protocols can accomplish a performance increase compared to unassisted protocols.
- Introducing entanglement resources transforms the capacity formula from multi-letter to single-letter form
- Dimension bound is an open problem — also for quantum wiretap channel, quantum broadcast channel, squashed entanglement, etc.

Thank you

Theorem

Let $|\omega_{ABK}\rangle, |\sigma_{SRG_1G_2}\rangle = |\Psi_{SR}\rangle \otimes |\Phi_{G_1G_2}\rangle$ in $\mathcal{H}_S^{\otimes 2} \otimes \mathcal{H}_G^{\otimes 2}$. Let $W_{SG_1 \rightarrow A^n}$ be a full-rank partial isometry, and denote $|\sigma_{A^nRG_2}\rangle = W_{SG_1 \rightarrow A^n} |\sigma_{SRG_1G_2}\rangle$. Define

$$\mathcal{T}_{A \rightarrow K}(\rho_A) = |\mathcal{H}_A| \text{Tr}_B [op_{A \rightarrow BK}(|\omega_{ABK}\rangle)(\rho_A)]$$

where $op_{A \rightarrow B}(|i_A\rangle \otimes |j_B\rangle) \equiv |j_B\rangle \langle i_A|$. Then,

$$\int_{\mathbb{U}_{A^n}} dU_{A^n} \left\| \mathcal{T}_{A \rightarrow K}^{\otimes n}(U_{A^n} \sigma_{A^nR}) - \omega_K \otimes \sigma_R \right\|_1 \leq \sqrt{\frac{|\mathcal{H}_S|}{|\mathcal{H}_G|} 2^{-nH(A|K)_\omega + n\varepsilon_n}}$$

$$\int_{\mathbb{U}_{A^n}} dU_{A^n} \left\| \mathcal{T}_{A \rightarrow K}^{\otimes n}(U_{A^n} \sigma_{A^nRG_2}) - \omega_K \otimes \sigma_{RG_2} \right\|_1 \leq \sqrt{|\mathcal{H}_S| |\mathcal{H}_G| 2^{-nH(A|K)_\omega + n\varepsilon_n}}$$

where the integral is over the Haar measure on all unitaries U_{A^n} .

Achievability Scheme

